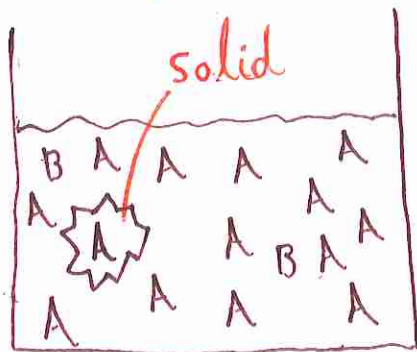


Ch. 25 \rightarrow solid-liquid solutions.

Freezing point depression: A = solvent
B = solute



Solid solvent is in equilibrium with the solvent in the solution.

$$\mu_A^*(s) = \mu_A(sol) = \mu_A^*(l) + RT \ln \alpha_A$$

Now, $\mu_A^*(s) - \mu_A^*(l) = -\Delta \bar{G}_{fusion}$ (pure substance)

$$\Rightarrow -\frac{\Delta \bar{G}_{fus}}{RT} = \ln \alpha_A \Rightarrow -\frac{1}{R} \left(\frac{\partial \Delta \bar{G}_{fus}/T}{\partial T} \right)_P = \left(\frac{\partial \ln \alpha_A}{\partial T} \right)_P$$

Recall (Raoult's Law) that $\alpha_A \rightarrow X_A$ as $X_A \rightarrow 1$.

$-\frac{\Delta \bar{H}_{fus}}{T^2}$ from Gibbs-Helmholtz

We now integrate from

T_{fus}^* (pure liquid, $X_A = 1 \Rightarrow \alpha_A = 1$) to T_{fus} (solution)

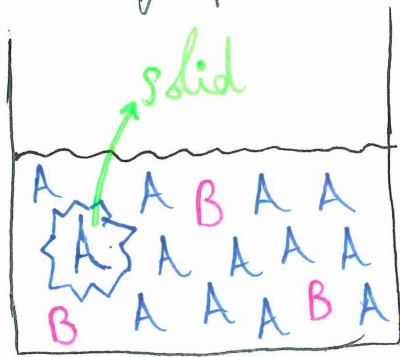
which gives $\ln \alpha_A \Big|_1^{\alpha_A} = \int_{T_{fus}^*}^{T_{fus}} \frac{\Delta \bar{H}_{fus}}{RT^2} dT = \ln \alpha_A$

Now, assume $\Delta \bar{H}_{fus}$ independent of T , and

put $\ln \alpha_A = \ln X_A = \ln(1 - X_B) = -X_B$ for X_B small

to get $-X_B = \frac{\Delta \bar{H}_{fus}}{R} \left(\frac{1}{T_{fus}^*} - \frac{1}{T_{fus}} \right) = \frac{\Delta \bar{H}_{fus}}{R} \cdot \frac{T_{fus} - T_{fus}^*}{T_{fus} T_{fus}^*}$

Freezing point depression, A = solvent, B = solute (68a)



- Solid is pure A because crystal doesn't support "B" defects.

- Component A in equilibrium at T_{fus}
 $\Rightarrow \mu_A^*(s) = \mu_A^{sol} = \mu_A^*(l) + RT_{fus} \ln X_A$
 using Raoult's Law

$$\Delta \bar{G}_{fus} \equiv \mu_A^*(l) - \mu_A^*(s) \text{ for the pure substance A.}$$

$$= 0 \text{ at } T_{fus}^* \Rightarrow \Delta \bar{S}_{fus} = \frac{\Delta \bar{H}_{fus}}{T_{fus}^*}$$

$$\text{Thus } \frac{\Delta \bar{G}_{fus}}{T_{fus}} = -R \ln X_A = \frac{\Delta \bar{H}_{fus}}{T_{fus}} - \Delta \bar{S}_{fus} = \Delta \bar{H}_{fus} \left(\frac{1}{T_{fus}} - \frac{1}{T_{fus}^*} \right)$$

Use a Taylor series: $\ln X_A = \ln(1 - X_B) = -X_B$ for X_B small

$$\text{to get } X_B = \frac{\Delta \bar{H}_{fus}}{R} \cdot \frac{(T_{fus}^* - T_{fus})}{T_{fus} T_{fus}^*}$$

Set $\Delta T_{fus} = T_{fus}^* - T_{fus}$ and approx. $T_{fus} T_{fus}^* = T_{fus}^{*2}$

to get $X_B = \frac{\Delta \bar{H}_{fus}}{R} \cdot \frac{\Delta T_{fus}}{T_{fus}^{*2}}$

Next change from mole fraction to molality, m.

$$m \equiv \frac{n_B}{1 \text{ kg of A}} = \frac{n_B}{1000 \text{ g of A}}$$

and $X_B \equiv \frac{n_B}{n_A + n_B} \approx \frac{n_B}{n_A} = \frac{m M_A}{1000}$ (assume n_B small)
 $M_A = \text{molar mass of A}$

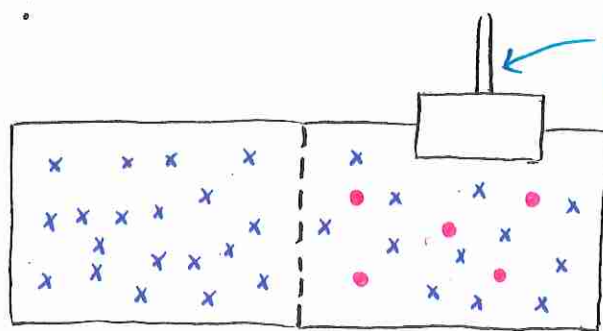
Finally, gets $\Delta T_{fus} = \left[\frac{M_A}{1000} \cdot \frac{R T_{fus}^{*2}}{\Delta \bar{H}_{fus}} \right] m = k_f m$

Similarly, $\Delta T_{vap} = k_b m$ with $k_b = \frac{M_A}{1000} \cdot \frac{R T_{vap}^{*2}}{\Delta \bar{H}_{vap}}$
for the boiling point elevation

Osmotic Pressure:

70

Usually the solvent is water.



x = solvent = 1
• = solute = 2

At equilibrium the chemical potential of water must be equal on the two sides.

↳ semipermeable membrane. won't let solute pass through.

On the right we have $\mu_i^{sd} = \mu_i^*(l) + RT \ln a_i$

For $x_i < 1$, $\mu_i^{sd} < \mu_i^*(l)$

Therefore somehow we must raise μ_i^{sd} .

We know $\left(\frac{\partial \mu_i}{\partial P}\right)_T = \bar{V}_i$ so we can increase the pressure on the piston in order to raise μ_i^{sd} .

Equil.: $\mu_i^*(l, P) = \mu_i^*(l, P + \pi) + RT \ln a_i$

$$\begin{aligned} \text{Now, } \mu_i^*(l, P + \pi) &= \mu_i^*(l, P) + \int_P^{P+\pi} \left(\frac{\partial \mu_i^*}{\partial P}\right)_T dP \\ &= \mu_i^*(l, P) + \pi \bar{V}_i^* \quad (\text{assume incompressible}) \end{aligned}$$

$$\Rightarrow \pi \bar{V}_i^* + RT \ln a_i = 0$$

put $\ln a_1 = \ln x_1 = \ln(1-x_2) \approx -x_2$ for x_2 small (71)

$$\text{Also put } x_2 = \frac{n_2}{n_1+n_2} \approx \frac{n_2}{n_1}$$

and put $n_1 \bar{V}_1^* \approx V$ (dilute solution)

$$\text{yielding } \pi \bar{V}_1^* = RT x_2 = RT \frac{n_2}{n_1}$$

$$\Rightarrow \pi V = n_2 RT$$

$$\Rightarrow \pi = CRT \quad C = \frac{n_2}{V} = \text{molarity of the solution}$$

ex 2.20 g of a polymer is dissolved in enough water to make 300 mL of solution. The osmotic pressure is found to be 7.45 torr at 20°C.

Determine the polymer's molar mass.

$$\pi = CRT = \frac{7.45 \text{ torr}}{760 \text{ torr atm}^{-1}} = 0.0098 \text{ atm}$$

$$C = \frac{\pi}{RT} = \frac{0.0098 \text{ atm}}{(0.082)(293 \text{ K})} = 4.07 \times 10^{-4} \text{ M}$$

In 1 L of solution we would have $2.20 \text{ g} \left(\frac{1000 \text{ mL}}{300 \text{ mL}} \right) = 7.33 \text{ g}$ of polymer.

Thus the molar mass is $\frac{7.33 \text{ g L}^{-1}}{4.07 \times 10^{-4} \text{ mol L}^{-1}} = 18000 \text{ g mol}^{-1} = 18 \text{ kDa}$.

For a strong electrolyte that dissociates into ν ions, we modify the colligative properties expressions:

$$\Delta T_{\text{fus}} = \nu k_f m \quad ; \quad \Delta T_{\text{vap}} = \nu k_b m \quad ; \quad \pi = \nu cRT$$

For a 2 molal aqueous solution of sucrose, to use $\pi = cRT$ we should compute c first:

$$c = \frac{2 \text{ mol sucrose}}{1000 \text{ g H}_2\text{O}} \cdot \frac{1000 \text{ g}}{1.003 \text{ L}} = 1.994 \text{ M} \quad \text{using density at 298 K of } 0.997 \text{ g/mL}$$

$$\pi = (1.994 \text{ M})(0.082)(298 \text{ K}) = 48.7 \text{ atm}$$

If we account for non-ideal behavior it turns out $\pi = 58.0 \text{ atm}$ is about 20% higher