

### Chem 3322 central force notes

Our goal is to solve the TISE for the central force problem where the potential is a function of  $r$  only.

$$-\frac{\hbar^2}{2m}\nabla^2\psi + V(r)\psi = E\psi \quad (1)$$

Try separation of variables:  $\psi(r, \theta, \phi) = R(r)P(\theta)Q(\phi)$

We need to transform the Laplacian into spherical polar coordinates, which is Eq. (5.49) in McQuarrie/Simon:

$$\nabla^2\psi = \frac{1}{r^2} \left[ \frac{\partial}{\partial r} \left( r^2 \frac{\partial \psi}{\partial r} \right) + \frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{\sin^2 \theta} \frac{\partial^2 \psi}{\partial \phi^2} \right] \quad (2)$$

The TISE is thus

$$-\frac{\hbar^2}{2m} \frac{1}{r^2} \left[ PQ \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{RQ}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) + \frac{RP}{\sin^2 \theta} \frac{d^2 Q}{d\phi^2} \right] + V(r)RPQ = ERPQ \quad (3)$$

Divide through by  $\psi = RPQ$  to get

$$-\frac{\hbar^2}{2m} \frac{1}{r^2} \left[ \frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{1}{P \sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) + \frac{1}{Q \sin^2 \theta} \frac{d^2 Q}{d\phi^2} \right] + V(r) = E \quad (4)$$

The expression in  $[\dots]$  is a function of  $r$  only. Why? Solve for it:

$$\left[ \frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{1}{P \sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) + \frac{1}{Q \sin^2 \theta} \frac{d^2 Q}{d\phi^2} \right] = -\frac{2mr^2(E - V(r))}{\hbar^2} \quad (5)$$

Therefore, the sum of the 2nd and 3rd terms in  $[\dots]$  must be a constant because they have angle dependence only. Call this constant  $-\lambda$ .

$$\frac{1}{P \sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) + \frac{1}{Q \sin^2 \theta} \frac{d^2 Q}{d\phi^2} = -\lambda \quad (6)$$

Rearrange Eq. (6) to obtain

$$\frac{1}{Q \sin^2 \theta} \frac{d^2 Q}{d\phi^2} = -\lambda - \frac{1}{P \sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) \quad (7)$$

and then again to obtain

$$\frac{1}{Q} \frac{d^2 Q}{d\phi^2} = \sin^2 \theta \left( -\lambda - \frac{1}{P \sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) \right) \quad (8)$$

Hence the left hand side of Eq. (8) is a function of  $\theta$  only, which can only hold if

$$\frac{1}{Q} \frac{d^2 Q}{d\phi^2} = -m^2 \quad (9)$$

We have achieved our goal! The three separated equations are:

$$\frac{d^2 Q}{d\phi^2} = -m^2 Q \quad (10)$$

$$\frac{1}{\sin \theta} \frac{d}{d\theta} \left( \sin \theta \frac{dP}{d\theta} \right) - \frac{m^2}{\sin^2 \theta} P = -\lambda P \quad (11)$$

$$-\frac{\hbar^2}{2m} \frac{1}{r^2} \left[ \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) - \lambda R \right] + V(r)R = ER \quad (12)$$

A brief table of some of the  $P_{\ell m}(\theta)$  (unnormalized)

$$\begin{aligned}
 \ell = 0 & & P_{00} &= 1 \\
 \ell = 1 & & P_{10} &= \cos \theta \\
 & & P_{11} &= \sin \theta \\
 \ell = 2 & & P_{20} &= 3 \cos^2 \theta - 1 \\
 & & P_{21} &= \sin \theta \cos \theta \\
 & & P_{22} &= \sin^2 \theta \\
 \ell = 3 & & P_{30} &= 5 \cos^3 \theta - 3 \cos \theta \\
 & & P_{31} &= \sin \theta (5 \cos^2 \theta - 1) \\
 & & P_{32} &= \sin^2 \theta \cos \theta \\
 & & P_{33} &= \sin^3 \theta
 \end{aligned}$$

Spherical harmonics, normalized with respect to the  $(\theta, \phi)$  parts of 3d integration in spherical polar coordinates:

$$\begin{aligned}
 \ell = 0 & & \mathcal{Y}_{00} &= \left(\frac{1}{4\pi}\right)^{1/2} \\
 \ell = 1 & & \mathcal{Y}_{10} &= \left(\frac{3}{4\pi}\right)^{1/2} \cos \theta \\
 & & \mathcal{Y}_{11} &= \left(\frac{3}{4\pi}\right)^{1/2} \sin \theta \begin{cases} \sin \phi \\ \cos \phi \end{cases} \\
 \ell = 2 & & \mathcal{Y}_{20} &= \left(\frac{5}{16\pi}\right)^{1/2} (3 \cos^2 \theta - 1) \\
 & & \mathcal{Y}_{21} &= \left(\frac{15}{4\pi}\right)^{1/2} \sin \theta \cos \theta \begin{cases} \sin \phi \\ \cos \phi \end{cases} \\
 & & \mathcal{Y}_{22} &= \left(\frac{15}{16\pi}\right)^{1/2} \sin^2 \theta \begin{cases} \sin 2\phi \\ \cos 2\phi \end{cases}
 \end{aligned}$$

Notice that, for each  $\ell$ , there are  $2\ell + 1$  different spherical harmonics. Using  $x = r \sin \theta \cos \phi$ ,  $y = r \sin \theta \sin \phi$ ,  $z = r \cos \theta$  and various trig. identities, the spherical harmonics up to  $\ell = 2$  can be rewritten as: (normalized)

$$\begin{array}{lll}
 \ell = 0 & \mathcal{Y}_{00} = \left(\frac{1}{4\pi}\right)^{1/2} & \text{“s” orbital} \\
 \ell = 1 & \mathcal{Y}_{10} = \left(\frac{3}{4\pi}\right)^{1/2} z/r & \text{“}p_z\text{” orbital} \\
 & \mathcal{Y}_{11}^{\sin} = \left(\frac{3}{4\pi}\right)^{1/2} y/r & \text{“}p_y\text{” orbital} \\
 & \mathcal{Y}_{11}^{\cos} = \left(\frac{3}{4\pi}\right)^{1/2} x/r & \text{“}p_x\text{” orbital} \\
 \ell = 2 & \mathcal{Y}_{20} = \left(\frac{5}{16\pi}\right)^{1/2} (3z^2 - r^2)/r^2 & \text{“}d_{z^2}\text{” orbital} \\
 & \mathcal{Y}_{21}^{\sin} = \left(\frac{15}{4\pi}\right)^{1/2} yz/r^2 & \text{“}d_{yz}\text{” orbital} \\
 & \mathcal{Y}_{21}^{\cos} = \left(\frac{15}{4\pi}\right)^{1/2} xz/r^2 & \text{“}d_{xz}\text{” orbital} \\
 & \mathcal{Y}_{22}^{\sin} = \left(\frac{15}{4\pi}\right)^{1/2} xy/r^2 & \text{“}d_{xy}\text{” orbital} \\
 & \mathcal{Y}_{22}^{\cos} = \left(\frac{15}{16\pi}\right)^{1/2} (x^2 - y^2)/r^2 & \text{“}d_{x^2-y^2}\text{” orbital}
 \end{array}$$

You can now understand the origin of the traditional designations of  $p_x$ ,  $p_y$ ,  $p_z$ , etc.